Exoplanets III. Outline

- Fundamentals of Stellar Physics
- Asteroseismology
- Finding of Invisible Companions in Binary Systems

Hiromoto Shibahashi Department of Astronomy

I. Fundamentals of Stellar Physics

Observational Facts

Characteristic quantities

- □ Size R_s
- □ Mass
- **Luminosity**
- Dyn timescale
- □ Therm timescale

 $R_{sun} = 7 \ 10^8 \text{ m}$ $M_{sun} = 2 \ 10^{30} \text{ kg}$ $L_{sun} = 4 \ 10^{26} \text{ W}$ $\tau_{dyn} = (GM/R^3)^{-1/2} \sim 1 \text{ hr}$ $\tau_{KH} = GM^2/(RL) \sim 10^7 \text{ yr}$

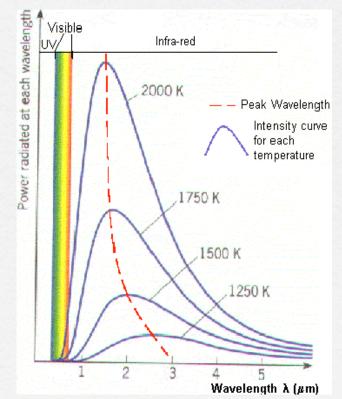
Light from a star

Thermal radiation from a body with T

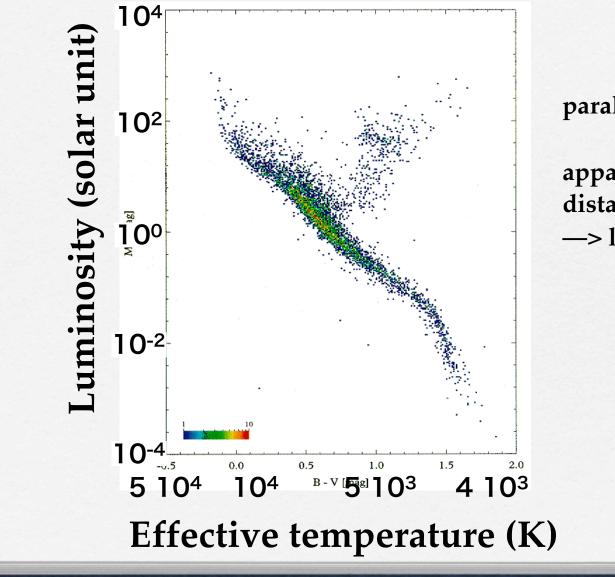
$$B_\lambda(T) = rac{2\pi hc^2}{\lambda^5} rac{1}{\exp(hc/\lambda kT)-1}$$

- Temperature determines spectrum
- Colour indicates temperature

 $\lambda_{max}T=2.9 \ 10^{-3} \ {\rm m} \ {\rm K}$



Effective Temperature vs Luminosity



parallax —> distance

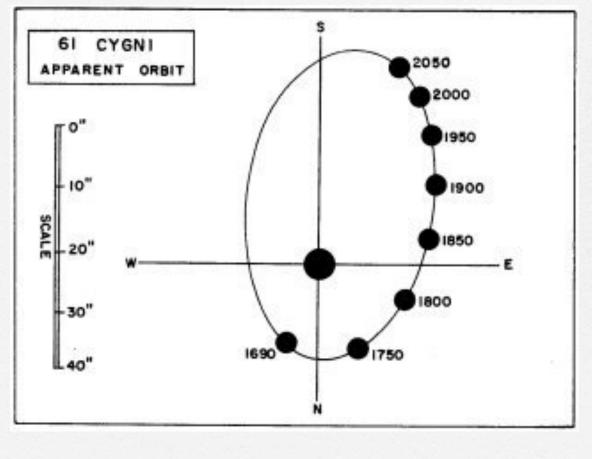
apparent brightness + distance —> luminosity

Determination of Stellar Mass

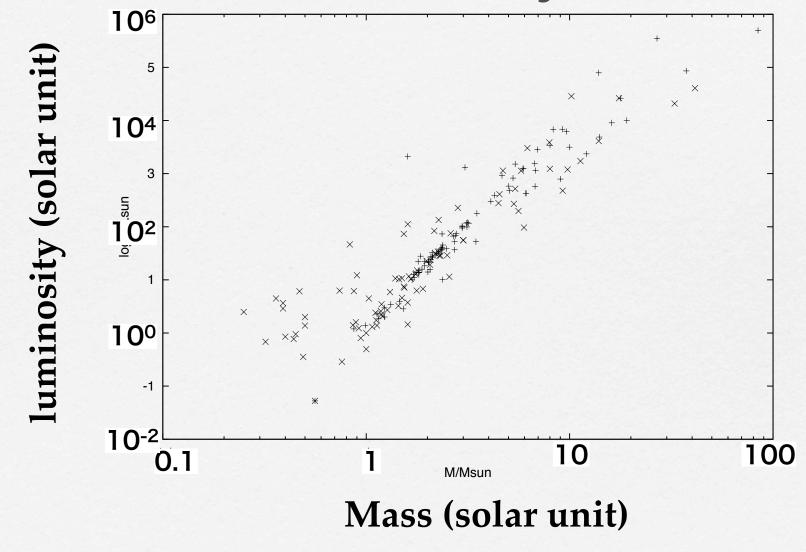
Observables:

(i) period *P* (ii) parallax *p* \longrightarrow distance *d* (iii) orbit size $r_1 \& r_2$

gravity = centrifugal force $2Gm_1m_2/(r_1 + r_2)^2$ $= 4\pi^2(m_1r_1 + m_2r_2)/P^2$ $r_1/r_2 = m_2/m_1$

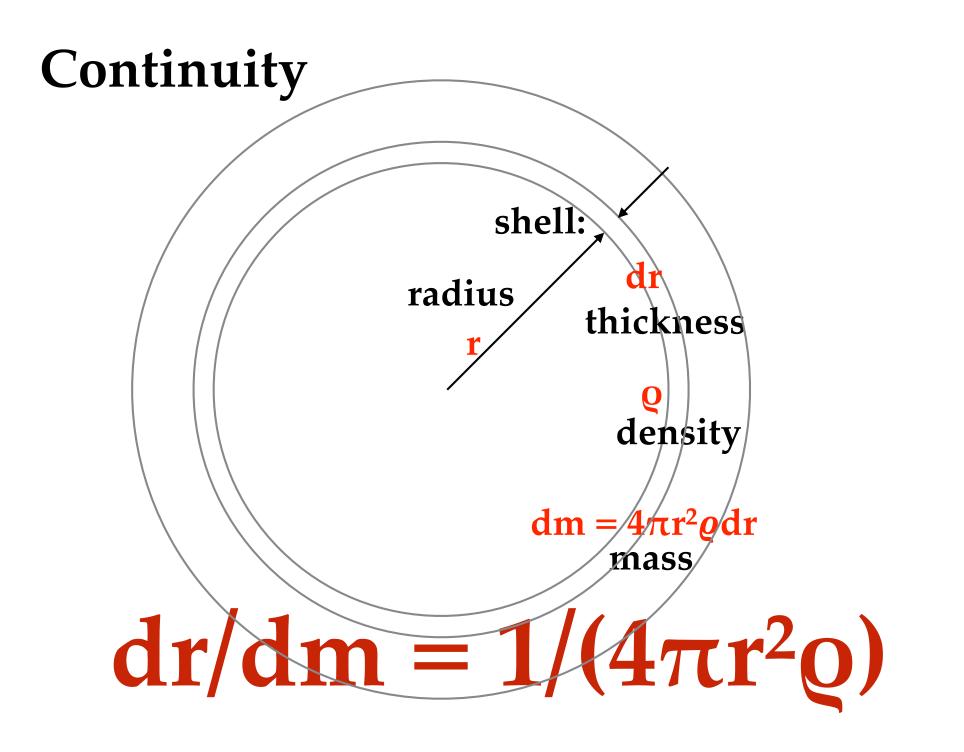


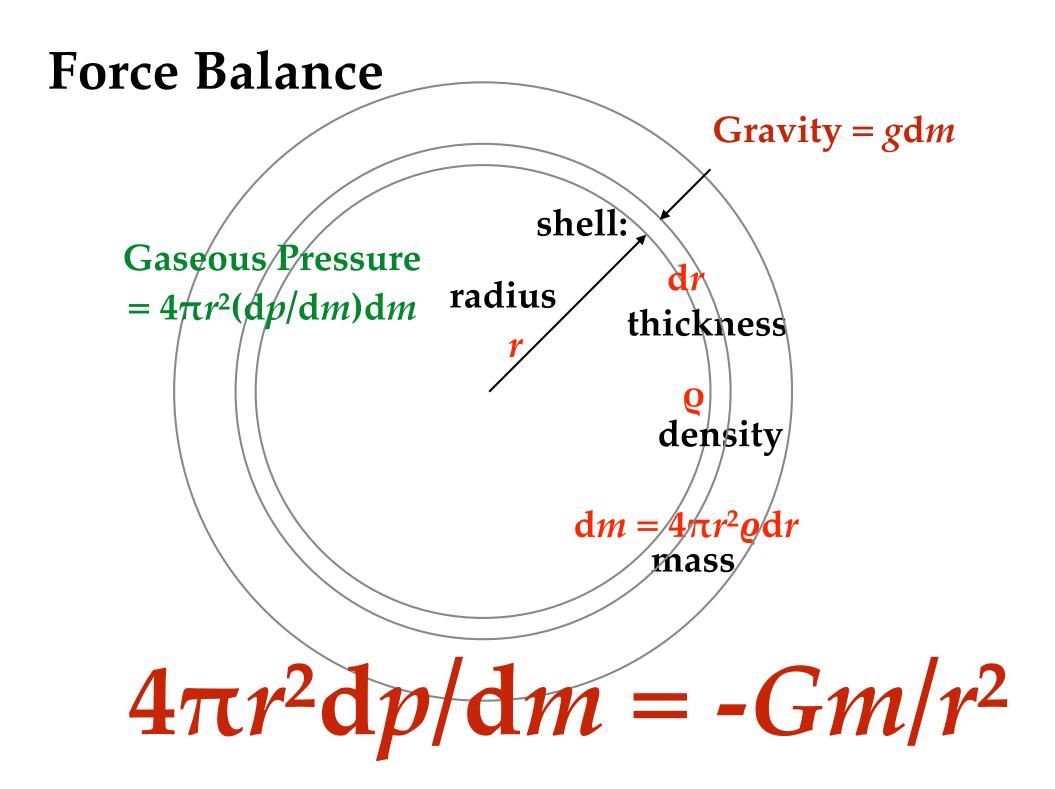
Mass-Luminosity relation

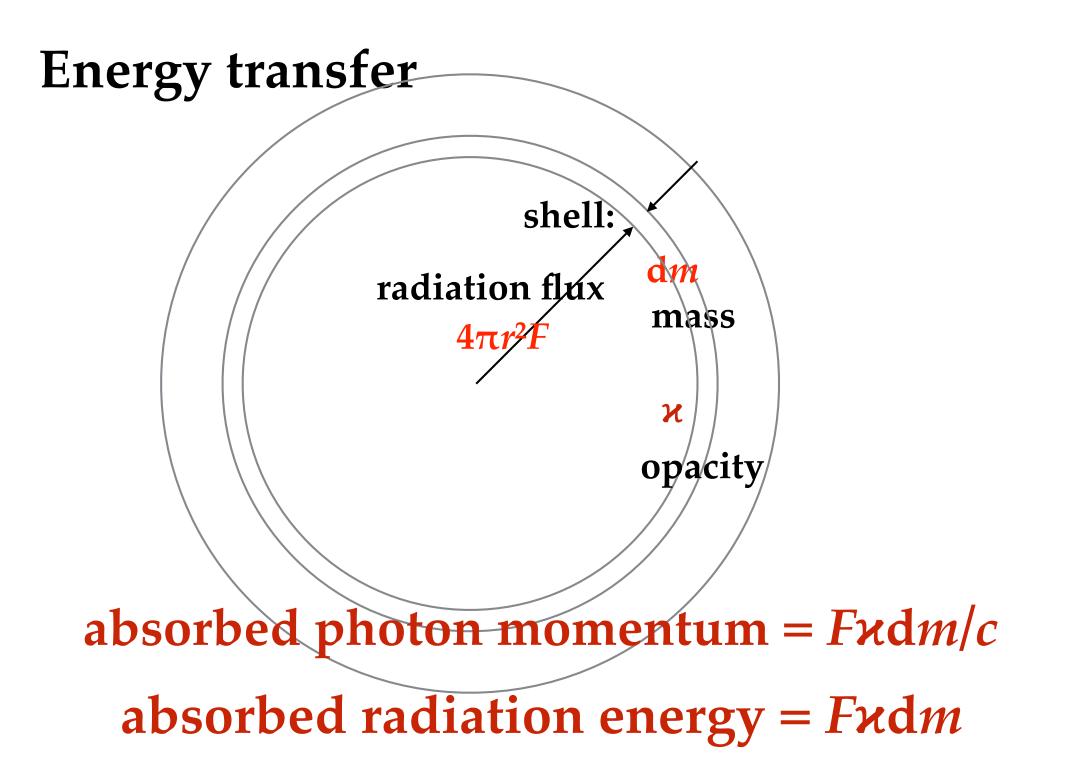


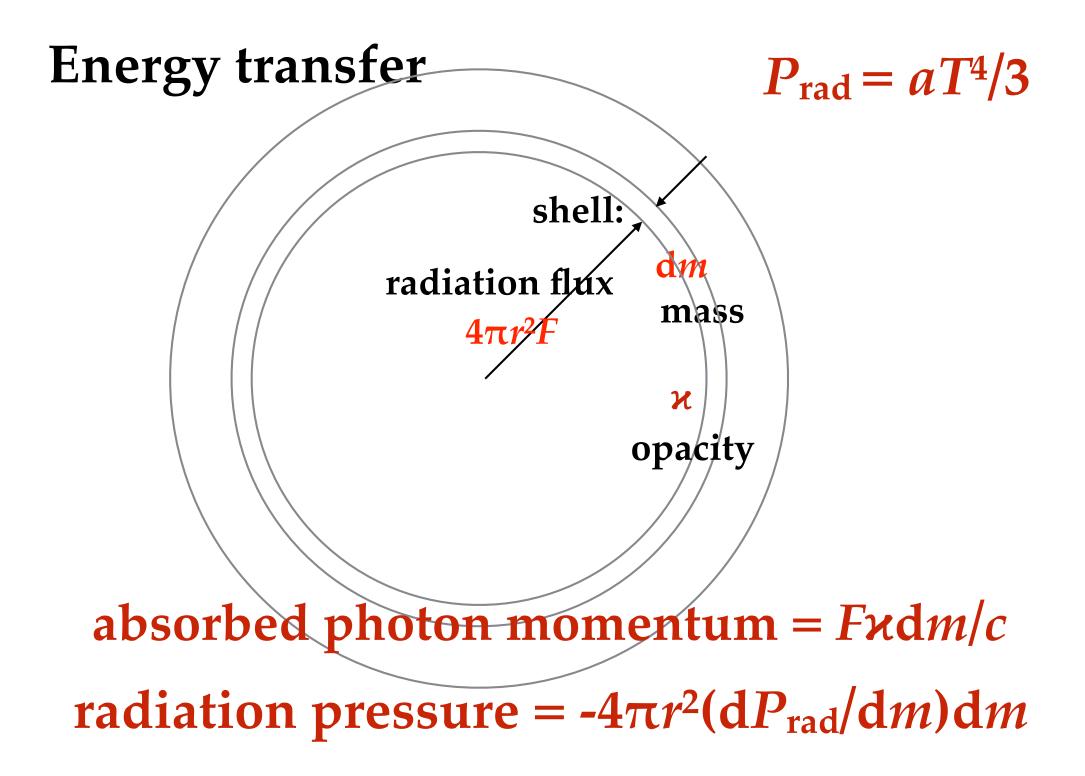
Theoretical consideration

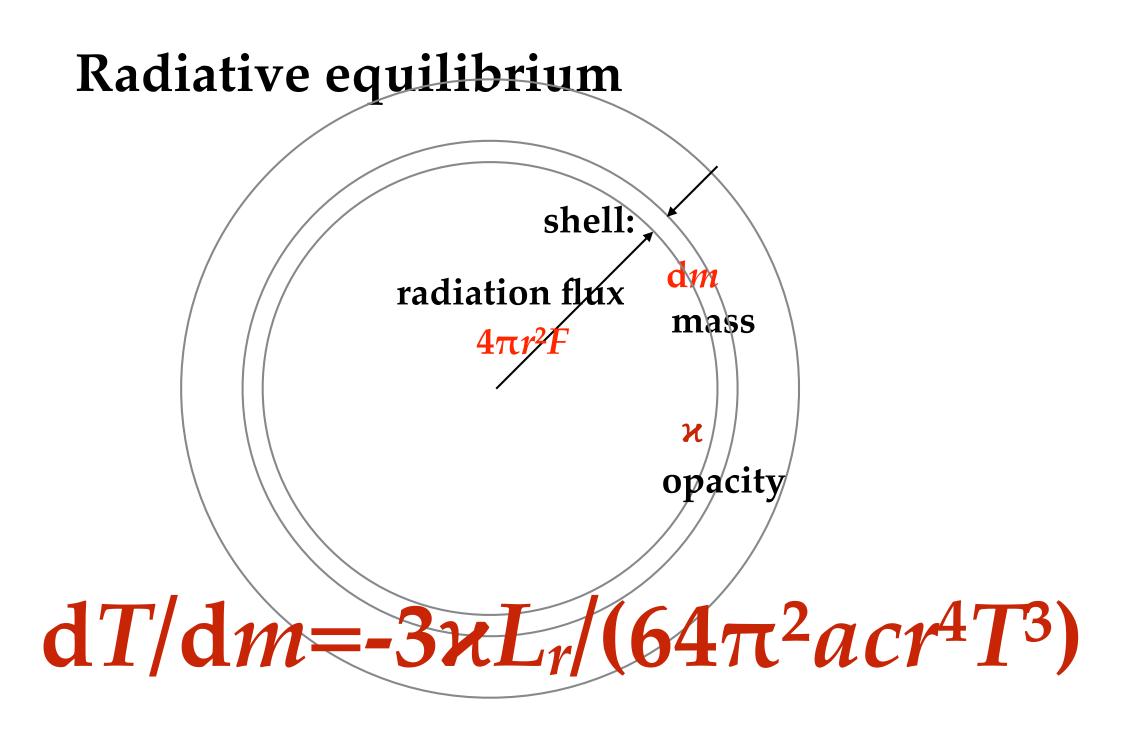
Why are stars shining?











Equilibrium state

 $dr/dm = 1/(4\pi r^2 \varrho)$ $dp/dm = - Gm\varrho/(4\pi r^4)$ $dT/dm = - 3\kappa L_r/(64\pi^2 a c r^4 T^3)$

Rough estimate: Differential Eq. --> Difference Eq. LHS: Difference between Surface and Center RHS: Averaged values

Rough estimate: Differential Eq. --> Difference Eq. LHS: Difference between Surface and Center RHS: Averaged values

 $\mathrm{d}r/\mathrm{d}m = 1/(4\pi r^2 \varrho)$

LHS $\approx R/M$ RHS $\approx (4\pi)^{-1}(R/2)^{-2}(\varrho_c/2)^{-1}$

 $: \mathbb{Q}_{c} \approx (2/\pi)(M/R^{3})$

Rough estimate: Differential Eq. --> Difference Eq. LHS: Difference between Surface and Center RHS: Averaged values

 $\mathrm{d}p/\mathrm{d}m = - Gm/(4\pi r^4)$

LHS $\approx -p_c/R$ RHS $\approx -G/(4\pi) (M/2)(R/2)^{-4}$

 $p_{\rm c} \approx (2/\pi)(GM^2/R^4)$

The central temperature

Ideal gas

$$p = nkT$$

 $= (\varrho/\mu m_u)kT$
 $\therefore T_c \approx (k/\mu m_u)^{-1}GM/R$
 $\approx 10^7 \text{ K for } M_{\text{sun}} R_{\text{sun}}$

n : particle numbers μ : mean molecular weight *k* = Boltzmann constant (1.38 10⁻²³ J/K) *m*_u = atomic weight (1.66 10⁻²⁵ kg)

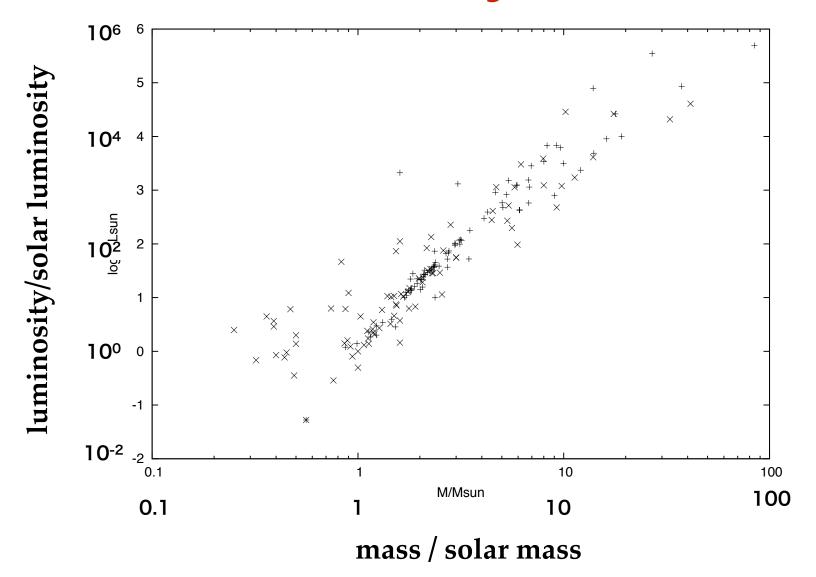
Rough estimate: Differential Eq. --> Difference Eq. LHS: Difference between Surface and Center RHS: Averaged values

 $dT/dm = -3\kappa L_r/(64\pi^2 a c T^3 r^4)$

LHS $\approx -T_c/M$ RHS $\approx -3 < \kappa > (L/2)/(64\pi^2 ac) (T_c/2)^{-3} (R/2)^{-4}$

: $L \approx \pi^2/(3 < \kappa >) \{acG^4/(k/m_u)^4\} \mu^{-4} M^3$

Mass-Luminosity relation



Radiation from a Star

Stefan-Boltzmann law: Radiation energy flux is proportional to *T*⁴

$$L = A \int B_{\lambda} d\lambda = A \sigma T_{\rm eff}^4$$

 $A = \text{surface area (m²)} = 4\pi R^2$ (*R*: stellar radius)

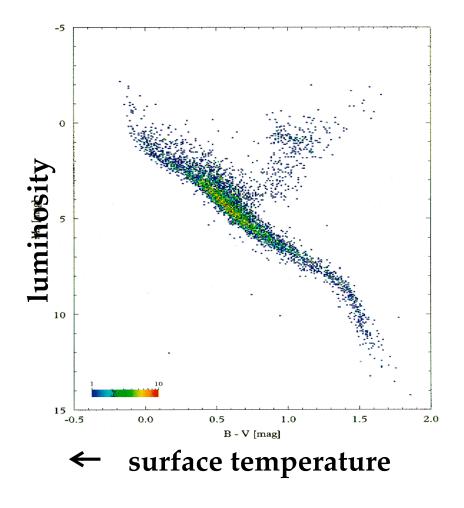
Main Sequence

$L \approx \pi^2/(3 < \kappa >) \{acG^4/(k/m_u)^4\}\mu^{-4}M^3$

Normalizing with the solar values,

 $<\!\!\kappa\!\!>\!\!L/(<\!\!\kappa\!\!>\!\!L)_{\rm sun}\approx (M/M_{\rm sun})^3$

Since $\sigma T_{eff}^4 = L/(4\pi R^2)$, $(T_{eff}/T_{eff,sun})^4 = (L/L_{sun})(R/R_{sun})^{-2}$ $\approx (L/L_{sun})(M/M_{sun})^{-2}$ $\approx (L/L_{sun})^{1/3}$ $\therefore L/L_{sun} \propto (T_{eff}/T_{eff,sun})^{12}$



Why are stars shining?

Nuclear fusion?

No!

Why are stars shining?

- Self gravity is supported by pressure.
- High gaseous pressure needs high temperature.
- Central temperature reaches 10⁷ K.
- Energy flows from hot to cool regions.

Why are stars shining?

Simply because stars are hot !

Energy flows from hot to cool region.

Break

Stellar Evolution

Star: Energy loosing system

Energy loss = Cooling

Cooling timescale = $\int c_v T dm/L$ $\approx 10^7$ yr for the Sun !

Lifetime $\propto M/L \propto M^{-2}$

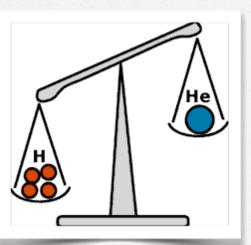
Necessity for sustaining mechanism

Nuclear fusion?

That's it!

Why can stars shine so long?

mass = energy



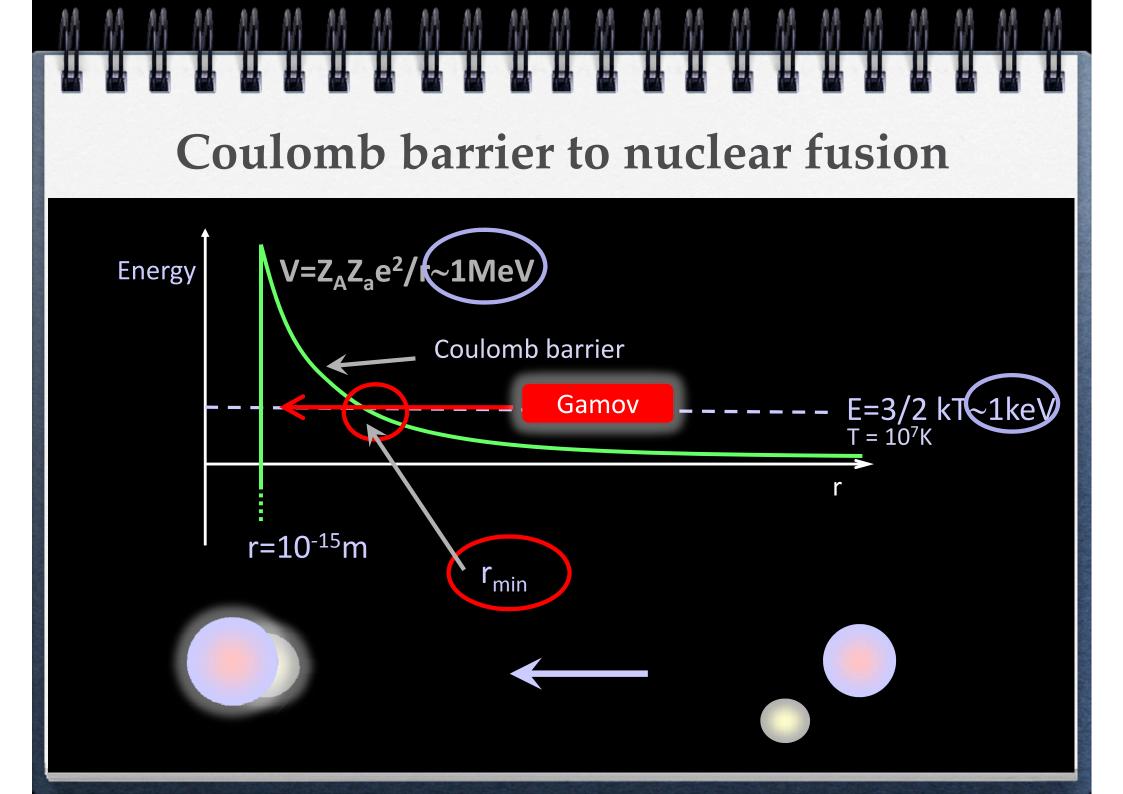
 H
 atomic weight
 1.008

 He
 atomic weight
 4.002

 {4m(H) - m(He)}/4 = 0.007

Assume 1. Solar composition: pure H 2. 10% of H converted to He

 $E_{nuc} \approx 0.007 \ (0.1 M_{sun})c^2$ $\approx 1.3 \ 10^{44} \ J$ $t_{nuc} = E_{nuclear} / L_{sun}$ $\approx 10^{10} \ yr$



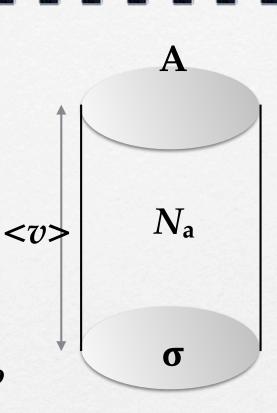
Nuclear reactions

 $r = N_A N_a < \sigma v >$

Maxwell-Boltzmann distribution for v

$$<\!\!\sigma v\!\!>=\int_0^\infty \sigma(v) v rac{1}{(2\pi m_\mu kT)^{3/2}} \mathrm{e}^{-E/kT} 4\pi p^2 \, dp$$

$$<\!\!\sigma v\!\!>=\int_0^\infty rac{8\pi}{m_\mu^{1/2}} rac{1}{(2\pi kT)^{3/2}} \sigma(E) E\,{
m e}^{-E/kT}\,dE$$



Cross section

Geometrical factor $\sigma = \pi \lambda^2 \approx \pi h^2 / (2m_\mu E)$ Penetration factor $P = \exp(-2\pi^2 r_{\min}/\lambda)$ $\propto \exp(-\beta E^{1/2})$

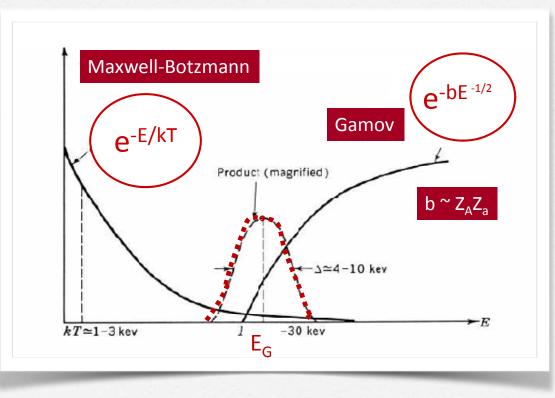
Nuclear factor S(E)



 $\sigma(E) = E^{-1} \exp(-\beta E^{1/2}) S(E)$

Nuclear reactions

$$<\!\!\sigma v\!\!>= \left(rac{8}{\pi m_{\mu}}
ight)^{1/2} \left(rac{1}{kT}
ight)^{3/2} \int_{0}^{\infty} S(E) \exp\left(-rac{E}{kT} - eta E^{-1/2}
ight) \, dE$$

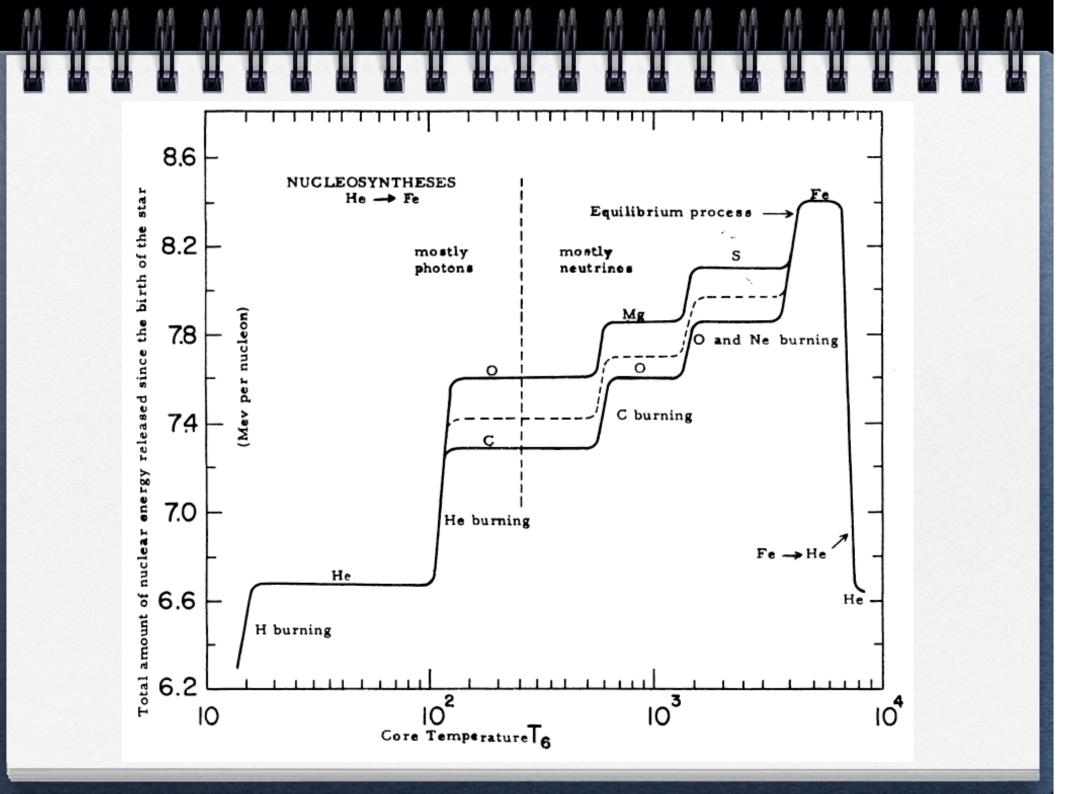


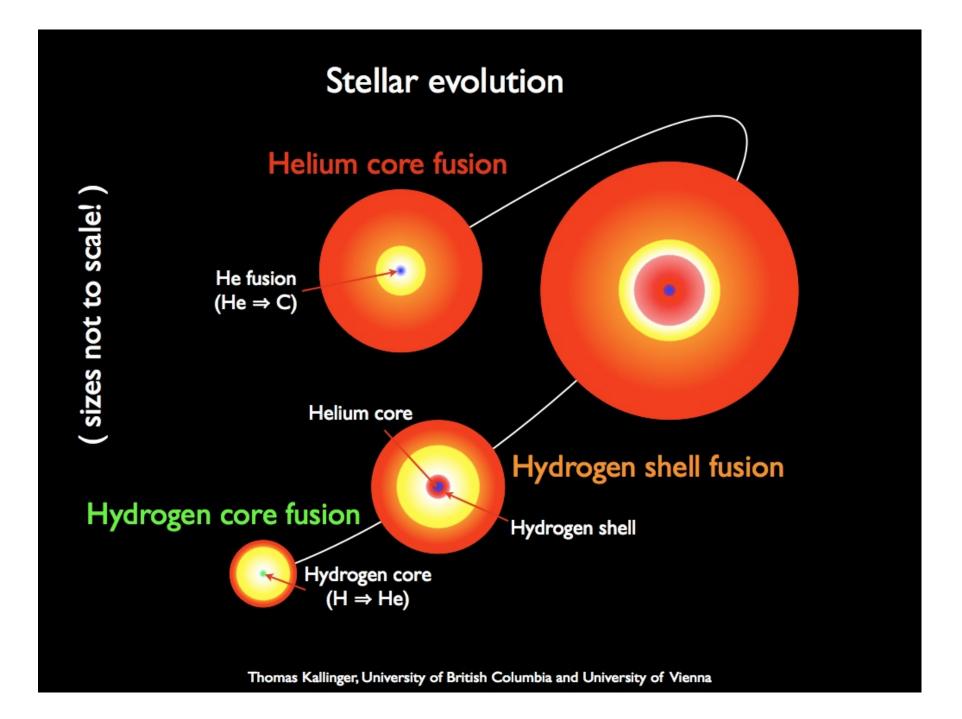
Gamow peak :
$$T \approx 10^7 \,\mathrm{K}$$

 $\approx T_{\rm c}$

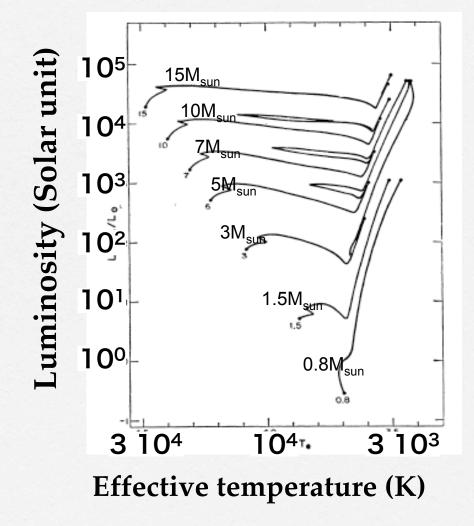
Is this a coincidence?

No! Stellar radius is adjusted so that $R \approx (k/\mu m_u)^{-1} GM/T_c$





Stellar Evolution



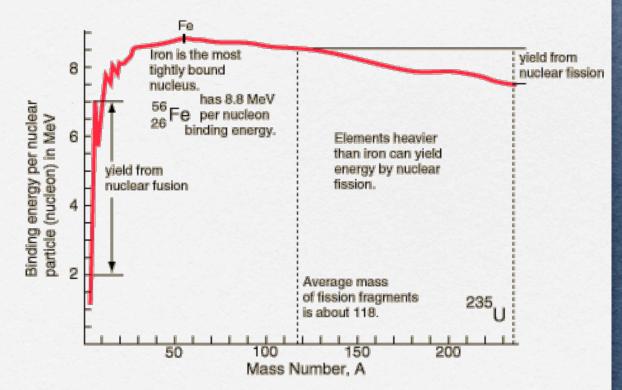
 $\Box \quad Mass-Luminosity \ relation \\ L \propto M^{\alpha}$

 \Box $\tau \propto M/L \propto M^{1-\alpha}$

 Masin Sequence -> Red Giants

Why Main Sequence?

Energy liberated by He and heavier nuclei is ~1/10 of the case of H burning



Hydrogen

Four hydrogen nuclei (4¹H) are converted to a helium nucleus (⁴He)

```
atomic weight of H = 1.008, so 4.032 by 4^{1}H atomic weight of He = 4.002
```

Hence, liberated energy per nucleus is proportional to (4.032 - 4.002)/4

Helium

Three helium nuclei (3⁴He) are converted to a nucleus of carbon (¹²C)

Atomic weight of He = 4.002, so 12.006 by 3^{4} He Atomic weight of C = 12.000

Hence, liberated energy per nucleus is proportional to (12.006 - 12.000)/12

Lifetime of He burning will be shorter than that of H burning by a factor of [(12.006 - 12.000)/12] / [(4.032 - 4.002)/4] **Essence of stellar evolution**

Toward gravitational contraction
 However, its timescale is not *GM*²/*RL* Residence by nuclear reactions
 Timescales are governed by nuclear reactions

Break

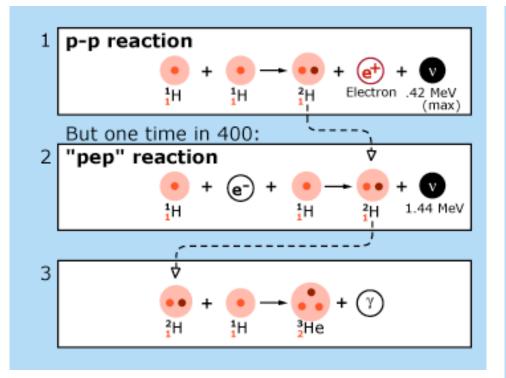
Is the Sun burning indeed?

In proof of nuclear reaction

- photon : diffusion process; 10 million
 years required to be transported from the center to the solar surface
- □ neutrino : no interaction with matter;

only 2 sec to reach the solar surface

 neutrino flux measurement is the only way to prove the nuclear reactions in the



 $1 \text{ MeV} = 1.6 \ 10^{-13} \text{ J}$

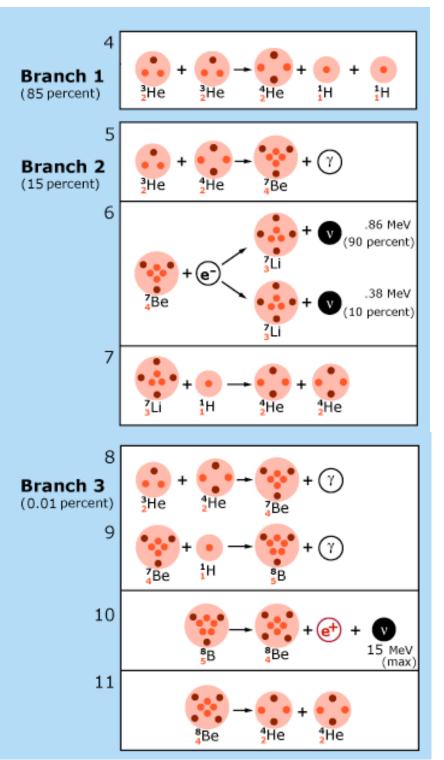
pp-chain

adopted from .

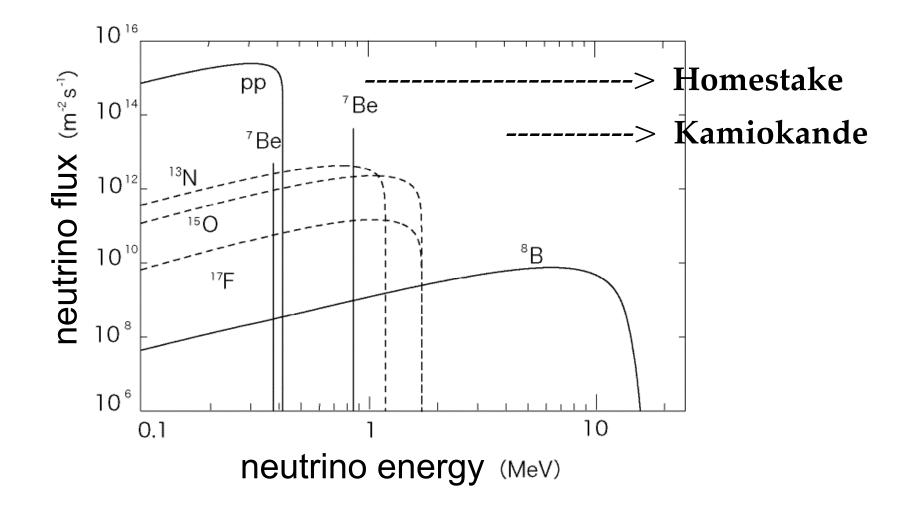
J.N. Bahcall,

Neutrinos from the Sun, Scientific American,

Volume 221, Number 1, July 1969, pp. 28-37.



Energy spectrum of solar neutrinos

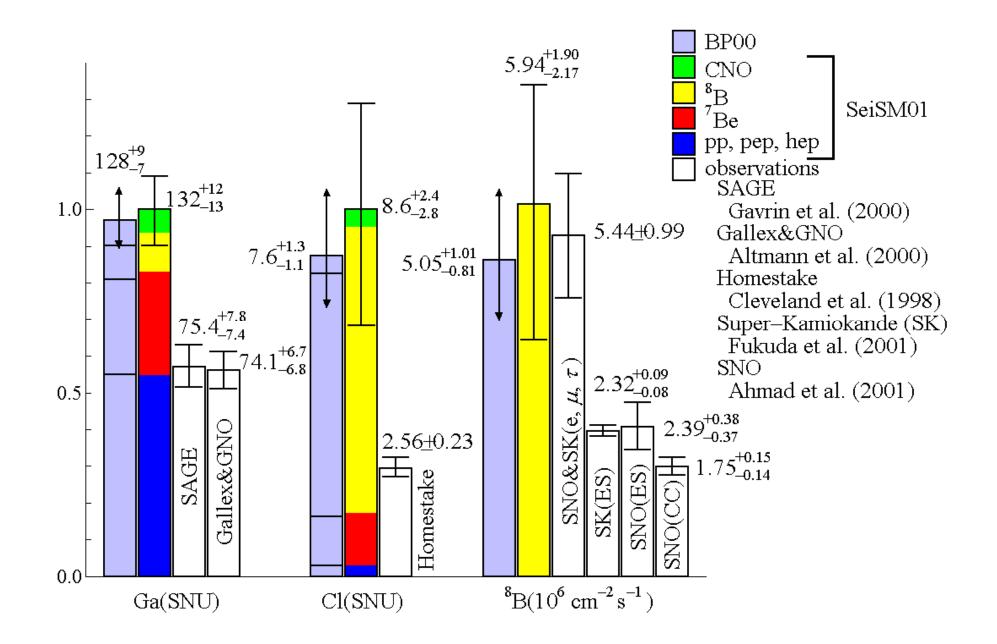


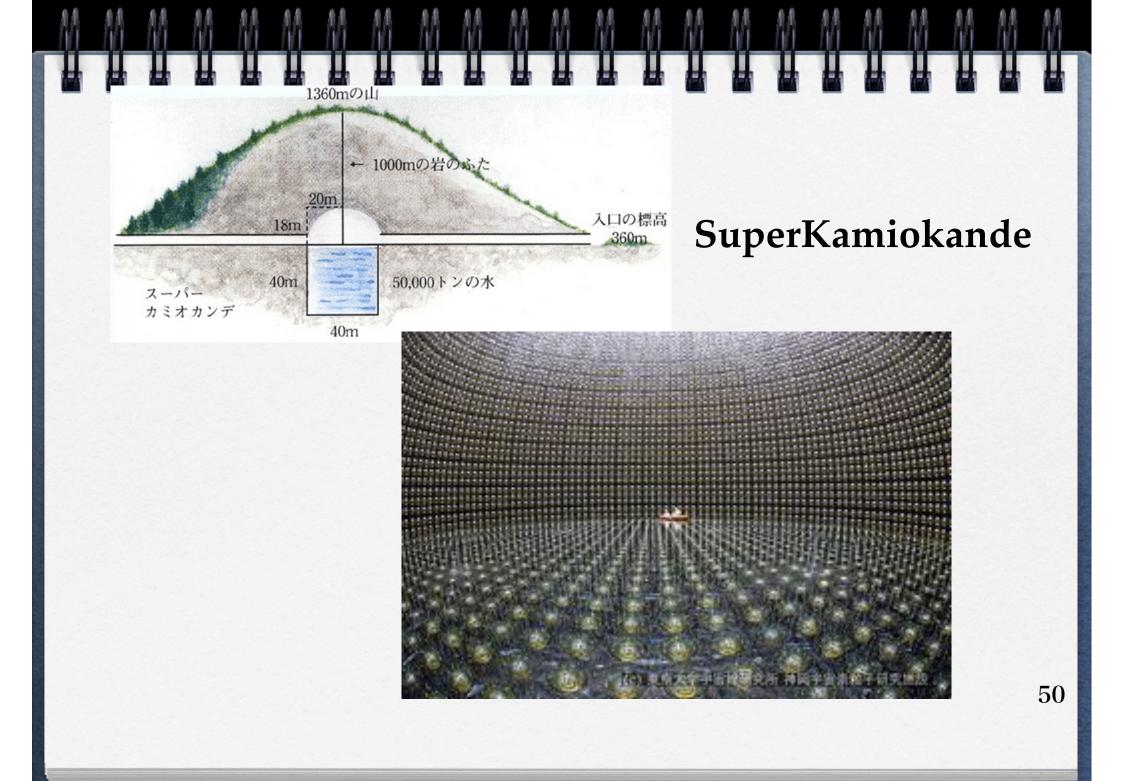
Solar neutrino problem

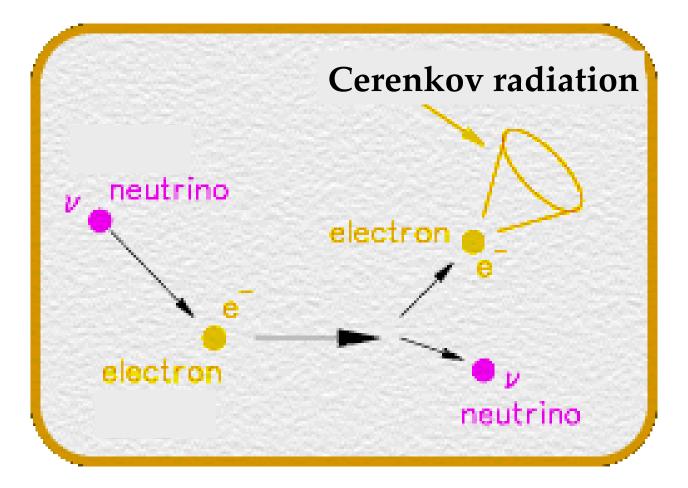
- pioneering work by R. Davis (1960-)
- Kamiokande (1987-)
- detected flux is about a half of theoretical prediction !
- experiments wrong ?
- solar models wrong ?
- particle physics wrong?



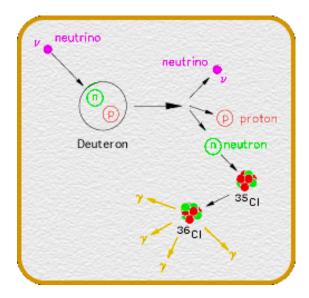
R. Davis & J. Bahcall



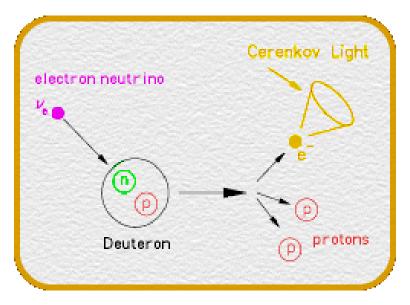




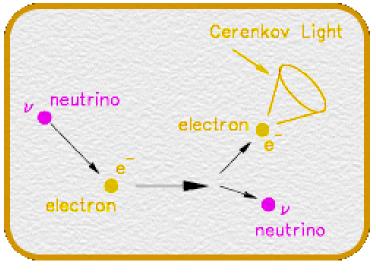
reaction occurring in SuperKamiokande



reaction II (NC mode)



reaction I (CC mode)



reaction III (ES mode)

Sudbury Neutrino Obs. experiment using D₂O

different sensitivity to neutrino types

SK (H₂O experiment)

- mainly sensitive to v_e
- but slightly sensitive also to $v_{\mu} \& v_{\tau}$

SNO (D₂O experiment)

• ES mode : the same as SK

experiment

- CC mode : only sensitive to v_e
- NC mode : sensitive to all of $v_{e}, v_{\mu} \& v_{\tau}$

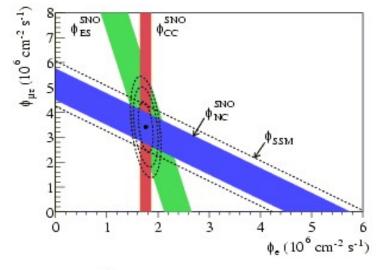
$$F_{SK} = F_{ve} + 0.15 (F_{v\mu} + F_{v\tau})$$
$$F^{CC}_{SNO} = F_{ve}$$

$$F^{\rm NC}_{\rm SNO} = F_{\rm ve} + (F_{\rm v\mu} + F_{\rm v\tau})$$

LHS : observables unknowns : F_{ve} & ($F_{v\mu}$ + $F_{v\tau}$)

solution to the solar neutrino problem: neutrino oscillations <-- non-zero mass of neutrinos!

FIG. 3: Flux of ⁸B solar neutrinos which are μ or τ flavor vs flux of electron neutrinos deduced from the three neutrino reactions in SNO. The diagonal bands show the total ⁸B flux as predicted by the SSM [11] (dashed lines) and that measured with the NC reaction in SNO (solid band). The intercepts of these bands with the axes represent the $\pm 1\sigma$ errors. The bands intersect at the fit values for ϕ_e and $\phi_{\mu\tau}$, indicating that the combined flux results are consistent with neutrino flavor transformation assuming no distortion in the ⁸B neutrino energy spectrum.



Solar neutrino problem

• Energy source = nuclear reactions

• 4H -> He + neutrinos time scale = several millions years

neutrino: no interaction

• photon: diffusion

End of Lecture I